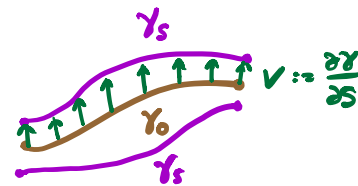


[Note: PS 4 due today, PS 5 posted and due on Apr 15.]

Recall:  $(M, g) \rightsquigarrow f(p, q) := \inf_{\gamma} L(\gamma)$  metric space structure

$$\left[ \frac{d}{ds} \Big|_{s=0} L(\gamma_s) = - \int_0^t \langle V, D_{\gamma} \gamma' \rangle dt + \langle V, \gamma' \rangle \Big|_{t=0} \right]$$



1st variation formula

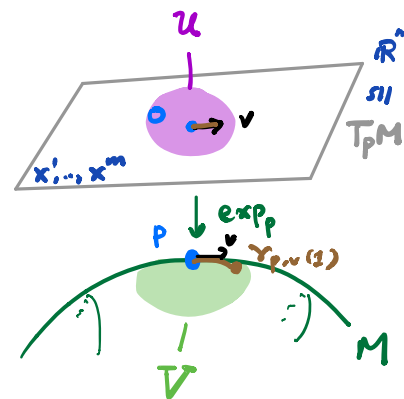
Geodesics  $\Leftrightarrow D_{\gamma} \gamma' \equiv 0 \Leftrightarrow$  critical pts to  $L$  (w.r.t. fixed end points)

Exponential map of  $(M, g)$  at  $p \in M$ :

$$\exp_p : \mathcal{U} \subseteq T_p M \xrightarrow[\cong]{\text{local diffeo.}} \mathcal{V} \subseteq M$$

$(\because d\exp_p(0) = \text{Id}_{T_p M})$

$$\exp_p(v) := \gamma_{p,v}(1)$$

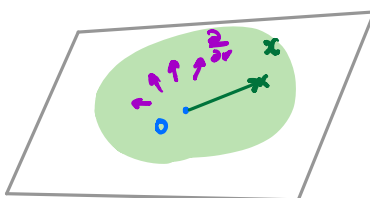


Eucl. coord.  $x^1, \dots, x^m$  on  $T_p M \cong \mathbb{R}^m$   $\xrightarrow{\exp_p}$  geodesic normal coord. near  $p \in M$ .

Properties of geodesic normal coord.  $x^1, \dots, x^m \rightsquigarrow g_{ij}, T_{ij}^k$

- (1)  $g_{ij}(0) = \delta_{ij}$  and  $T_{ij}^k(0) = 0$ . [In geod. normal coord.,  $g_{ij}(x) = \delta_{ij} + \underbrace{O(|x|^3)}_{\text{Curvature terms}}$ ]
- (2) Gauss Lemma:  $\sum_{j=1}^m g_{ij}(x) x^j = x^i$  — (\*)

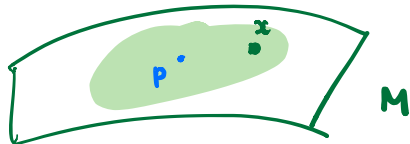
Geometric interpretation of (\*):



$$\begin{aligned} (T_p M, g_{ij}) &\rightsquigarrow \|x\|_g := [g(x,x)]^{1/2} \\ (\mathbb{R}^m, \delta_{ij}) &\rightsquigarrow r := |x| := [\delta(x,x)]^{1/2} = \left( \sum_{i=1}^m (x^i)^2 \right)^{1/2} \end{aligned}$$

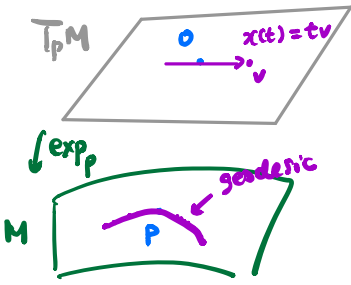
$$(*) \Leftrightarrow \|dr\|_g \equiv 1$$

i.e.  $\frac{\partial}{\partial r}$  is a unit vector field w.r.t.  $g$



$$\left[ \because dr = \sum_i \frac{x^i}{|x|} dx^i ; \|dr\|_g^2 = \sum_{i,j} g^{ij}(x) \frac{x^i}{|x|} \frac{x^j}{|x|} = \frac{|x|^2}{|x|^2} = 1 \right]$$

"Proof": (1)  $d\exp_p(0) = Id \Rightarrow g_{ij}(0) = \delta_{ij}$ .



$t \mapsto \exp_p(tv)$  is a geodesic on  $M \quad \forall v \in T_p M$

geodesic eq<sup>n</sup>  $\Rightarrow \frac{d^2 x^k}{dt^2} + \Gamma_{ij}^k(x(t)) \frac{dx^i}{dt} \frac{dx^j}{dt} = 0$   
 $x(t) = tv$

At  $t=0 \Rightarrow \Gamma_{ij}^k(0) v^i v^j = 0 \quad \forall v \in T_p M$

$\Rightarrow \Gamma_{ij}^k(0) = 0$  for all  $i, j, k$ .

(2)

Fix  $x_0 \in T_p M$ , consider  $\gamma_s : [0, 1] \rightarrow M$  given in g.n.c. by

$$\gamma_s(t) = t x_s$$

where  $s \mapsto x_s$  is a smooth curve on the

"sphere"  $\{x \in T_p M : |x| = |x_0|\} \subseteq T_p M$ .

Note: • For every fixed  $s$ ,

(\*)  $\rightarrow \gamma_s$  is a geodesic in  $M$ , with length  $|x_0|$  and speed  $= |x_0|$ .

(\*\*)  $\left\{ \begin{array}{l} \bullet \gamma_s(0) \equiv 0 \quad \forall s \\ \bullet \frac{d}{ds} \Big|_{s=0} \gamma_s(1) = \frac{d}{ds} \Big|_{s=0} x_s = v \end{array} \right.$

Notice:  $v \perp x_0$  w.r.t. flat metric

$$\text{i.e. } \sum_{i=1}^m v^i x_0^i = 0$$

Idea: Apply 1<sup>st</sup> var. formula for  $L$  to  $s \mapsto \gamma_s$ .

$$0 \stackrel{(*)}{=} \frac{d}{ds} \Big|_{s=0} L(\gamma_s) = - \int_0^1 \langle v, \underbrace{D_{x_0} x_0'}_0 \rangle_g dt + \langle v, \underbrace{\gamma_0'}_0 \rangle_g \Big|_{t=0}^{t=1}$$

$\stackrel{(*)}{=} 0$  by (\*)      by (\*\*),  $= \langle v, \gamma_0'(1) \rangle$

$$\Rightarrow \sum_{i,j} g_{ij}(x_0) x_0^i v^j = 0 \quad \forall v \text{ s.t. } v \perp^s x_0$$

$$\Rightarrow \sum_i g_{ij}(x_0) x_0^i \frac{\partial}{\partial x^j} \parallel x_0^j \frac{\partial}{\partial x^j} \quad \text{w.r.t. flat metric } \delta$$

Also,  $\sum_{i,j} g_{ij}(x_0) x_0^i x_0^j = \|\gamma_0'(1)\|_g^2 \stackrel{(*)}{=} |x_0|^2 = \sum_i (x_0^i)^2 \Rightarrow \sum_i g_{ij}(x_0) x_0^i = x_0^j$

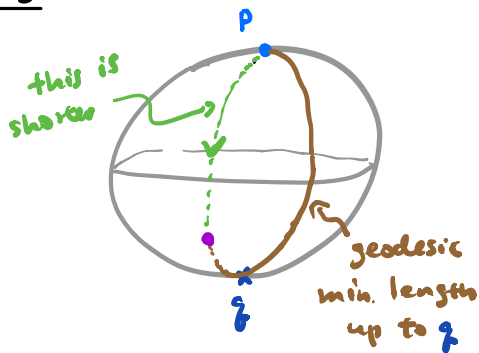
# Minimizing Properties of geodesics

Recall: "geodesics"  $\Leftrightarrow$  "straight lines" in  $(M, g)$

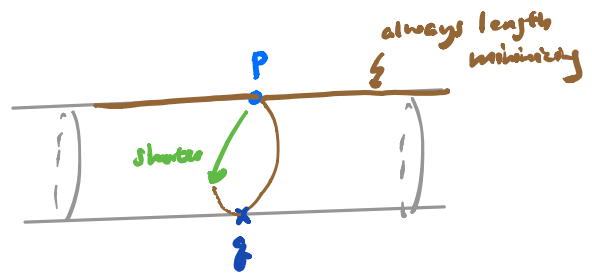
Q: How much does geodesics min. distance between 2 points?

FACT: "Short" geodesics minimizes length between its end points.  
 BUT not nec. true for "long" geodesics.

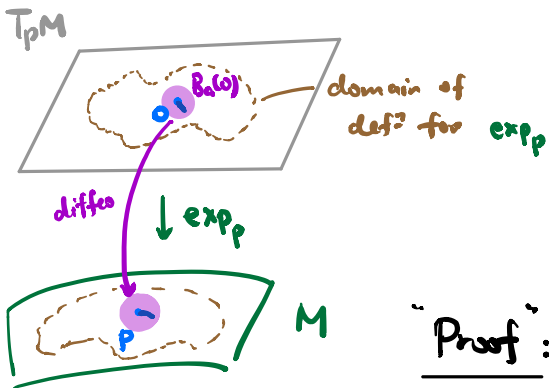
E.g.)  $M = S^2 \subseteq \mathbb{R}^3$



$M = S^1 \times \mathbb{R} \subseteq \mathbb{R}^3$



Prop: If  $\exp_p$  is a diffeo. onto its image on  $B_a(0) = \{x \in T_p M \mid |x| < a\}$ ,



THEN:  $\forall x \in B_a(0)$ , the ray  $x(t) = tx, 0 \leq t \leq 1$  is the unique length-minimizing geodesic from  $p$  to  $x(1)$ . (ie  $L(x(t)) = f(p, x(1))$ )

Proof: By Gauss Lemma,  $\|dx\|_g \equiv 1$ .

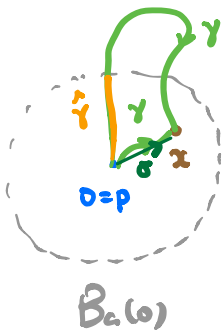
Case 1:  $\gamma$  stays inside  $B_a(0)$

$$L(\gamma) = \int ds \geq \int dr = |x| = L(\sigma) < a$$

$\uparrow$   
arc length  $\gamma$

Case 2:  $\gamma$  leaves the ball  $B_a(0)$  somewhere

$$L(\gamma) \geq L(\hat{\gamma}) \geq a$$



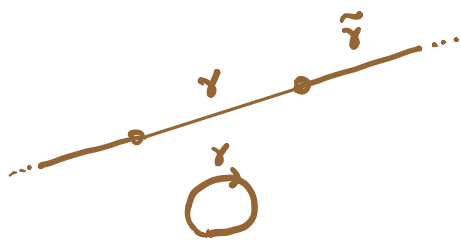
# Complete Riem. manifold

Recall:  $(M, g) \rightsquigarrow (M, \rho)$  metric space

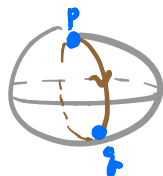
Metric space completeness: Every Cauchy seq. is convergent.

Q: How does this relate to  $g$ ?

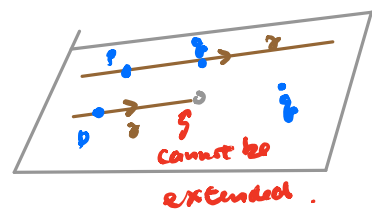
Def:  $(M, g)$  **geodesically complete** if any geodesic in  $M$  can be extended indefinitely, i.e.  $\exp_p: T_p M \rightarrow M$  is defined on the whole  $T_p M$ .  $\forall p \in M$



E.g.) Cpt Riem. mfd are always geod. complete



Non-E.g.:  $\mathbb{R}^2 \setminus \{0\} = M$  is not complete.



Hopf - Rinow Thm: Geodesic completeness  $\Leftrightarrow$  metric completeness.

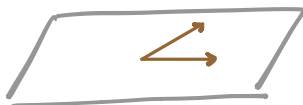
"Proof": HW exercise! [Key fact: complete  $\Rightarrow \exists$  min. geod.  $\gamma$  joining any two points]

Q: How does the "curvatures" of  $(M, g)$  affect the behavior of geodesics?

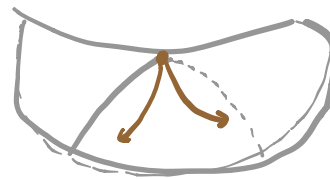
Recall:  $\Sigma^2 \subseteq \mathbb{R}^3$  surface



$K > 0$



$K = 0$



$K < 0$

Idea: In higher dimensions, the sectional curvatures affects the geodesics through 2<sup>nd</sup> variation of Length or "energy"

Def<sup>n</sup>: Given a curve  $\gamma: [a, b] \rightarrow (M, g)$ , the energy of  $\gamma$  is:

$$E(\gamma) := \int_a^b \|\gamma'(t)\|_g^2 dt$$

Remarks: (1)  $L(\gamma) := \int_a^b \|\gamma'(t)\|_g dt$  is indep of parametrization

BUT:  $E(\gamma)$  depends on the parametrization as well.

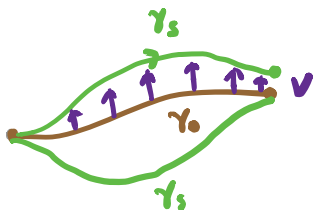
(2)  $E(\gamma)$  does not involve  $\sqrt{\cdot}$ , so it's easier to compute its derivative.

Note: (a) Hölder  $\Rightarrow L(\gamma)^2 \leq (b-a) \cdot E(\gamma)$ .

(b)  $f(p, q) := \inf_{\gamma} L(\gamma) = \inf_{\gamma} E(\gamma)^{1/2}$  (Take  $b-a = 1$ ).

1<sup>st</sup> & 2<sup>nd</sup> variations of energy

Setup: Let  $\gamma_s(t): [0, 1] \rightarrow M$ ,  $s \in (-\varepsilon, \varepsilon)$ , be a 1-parameter family of curves in  $M$ .



$$V(t) := \left. \frac{\partial}{\partial s} \right|_{s=0} \gamma_s(t)$$

variation field  
along  $\gamma_0$

(note:  $V \in T(\gamma_0^*TM)$ )

$$\left. \frac{1}{2} \frac{d}{ds} \right|_{s=0} E(\gamma_s) = - \int_0^1 \langle V, D_{\frac{\partial}{\partial t}} \gamma_0' \rangle_g dt + \langle \gamma_0', V \rangle \Big|_{t=0}^{t=1}$$

1<sup>st</sup> var. formula for energy Proof: Ex!

Remark: (i) We do NOT need to assume  $\gamma_0$  is p.b.a.l.

wt. fixed end pts  $\left\{ \begin{array}{l} \text{(ii) } \gamma \text{ is a critical pt. of } E \Leftrightarrow D_{\gamma} \gamma' \equiv 0, \text{ i.e. } \gamma \text{ geodesic} \\ \gamma \text{ is a critical pt. of } L \Leftrightarrow (D_{\gamma} \gamma')^{\perp} \equiv 0 \end{array} \right.$

Suppose:  $\gamma = \gamma_0$  is a closed geodesic.

Then, we have the 2<sup>nd</sup> variation formula for energy:

$$\frac{1}{2} \frac{d^2}{ds^2} \Big|_{s=0} E(\gamma_s) = \int_0^1 \left( \underbrace{\|D_{\frac{\partial}{\partial t}} V\|_g^2}_{\text{non-negative}} - \underbrace{R(\gamma', V, \gamma', V)}_{\text{sectional curvature term}} \right) dt$$

Proof: Write  $F(t, s) = \gamma_s(t)$ .

1<sup>st</sup> var. formula  $\forall s \in (-\epsilon, \epsilon)$  :  $\frac{1}{2} \frac{d}{ds} E(\gamma_s) = - \int_0^1 \langle D_{\frac{\partial}{\partial t}} \frac{\partial F}{\partial t}, \frac{\partial F}{\partial s} \rangle dt$   $V_s$

I.b.p.  $= \int_0^1 \langle \frac{\partial F}{\partial t}, D_{\frac{\partial}{\partial t}} \frac{\partial F}{\partial s} \rangle dt$

Differentiate in  $s$  again, and take  $s=0$ :

$$\begin{aligned} \frac{1}{2} \frac{d^2}{ds^2} \Big|_{s=0} E(\gamma_s) &= \int_0^1 \left[ \langle D_{\frac{\partial}{\partial s}} \frac{\partial F}{\partial t}, D_{\frac{\partial}{\partial t}} \frac{\partial F}{\partial s} \rangle + \langle \frac{\partial F}{\partial t}, D_{\frac{\partial}{\partial s}} D_{\frac{\partial}{\partial t}} \frac{\partial F}{\partial s} \rangle \right] \Big|_{s=0} dt \\ &= \int_0^1 \left[ \langle D_{\frac{\partial}{\partial s}} \frac{\partial F}{\partial t}, D_{\frac{\partial}{\partial t}} \frac{\partial F}{\partial s} \rangle + \langle \frac{\partial F}{\partial t}, D_{\frac{\partial}{\partial s}} D_{\frac{\partial}{\partial t}} \frac{\partial F}{\partial s} \rangle \right] \Big|_{s=0} dt \\ &\quad \left( \begin{array}{l} \text{torsion-free} \\ \downarrow \\ D_{\frac{\partial}{\partial s}} \frac{\partial F}{\partial t} = D_{\frac{\partial}{\partial t}} \frac{\partial F}{\partial s} \end{array} \right) \\ &\quad \left( \begin{array}{l} \text{using def.} \\ + R \end{array} \right) \rightarrow = \int_0^1 \|D_{\frac{\partial}{\partial t}} V\|_g^2 + \langle \frac{\partial F}{\partial t}, D_{\frac{\partial}{\partial t}} D_{\frac{\partial}{\partial s}} \frac{\partial F}{\partial s} \rangle \\ &\quad - R \left( \frac{\partial F}{\partial s}, \frac{\partial F}{\partial t}, \frac{\partial F}{\partial s}, \frac{\partial F}{\partial t} \right) \Big|_{s=0} dt \\ &= \int_0^1 \|D_{\frac{\partial}{\partial t}} V\|_g^2 - R(V, \gamma', V, \gamma') dt \\ &\quad - \int_0^1 \langle D_{\frac{\partial}{\partial t}} \frac{\partial F}{\partial t}, D_{\frac{\partial}{\partial s}} \frac{\partial F}{\partial s} \rangle \Big|_{s=0} dt \\ &\quad \underbrace{\hspace{10em}}_{=0 \because \text{no geodesic}} \end{aligned}$$

Def<sup>n</sup>:  $I(V, V) := \int_0^1 (\|D_{\frac{\partial}{\partial t}} V\|_g^2 - R(\gamma', V, \gamma', V)) dt \quad \forall V \in T(\gamma^* TM)$

is called the **index form** of a geodesic  $\gamma$ .

Note: polarize  $\leadsto I(V, W)$  symmetric bilinear.

Prop:  $V \in \ker I$  (ie.  $I(V, W) = 0 \quad \forall W \in T(\gamma^*TM)$ )

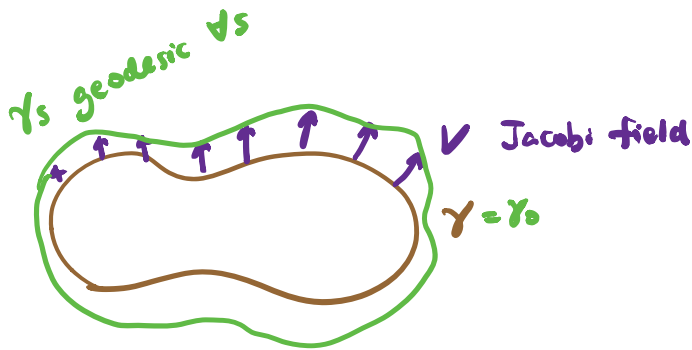
$\Leftrightarrow$

$$\boxed{D_{\frac{\partial}{\partial t}} D_{\frac{\partial}{\partial t}} V + R(\gamma', V) \gamma' = 0}$$

Jacobi field equation  
(2<sup>nd</sup> order linear).

Thm: Jacobi fields  $V$  arise exactly from 1-parameter family of geodesics  $\{\gamma_s\}_{s \in (-\epsilon, \epsilon)}$  s.t.  $\gamma_0 = \gamma$ .

Proof: Exercise HW.



Idea: geodesics  $\Leftrightarrow D_{\gamma'} \gamma' \equiv 0$   
(2<sup>nd</sup> order non-linear)

Jacobi field  $\Leftrightarrow$  "linearization"  
of  $D_{\gamma'} \gamma' \equiv 0$   
(2<sup>nd</sup> order linear)

$$D_{\gamma'_s} \gamma'_s \equiv 0 \quad \forall s \in (-\epsilon, \epsilon).$$

$$\Rightarrow \left. \frac{d}{ds} \right|_{s=0} (D_{\gamma'_s} \gamma'_s \equiv 0) \rightsquigarrow \text{Jacobi field eq.}^2$$